

Quantum tunneling in scintillating materials

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Goal:

Improvement of scintillator performance

**understanding the physical processes
involved in the scintillation mechanism**

**losses of fast scintillation light due to the
presence of traps in the host materials**

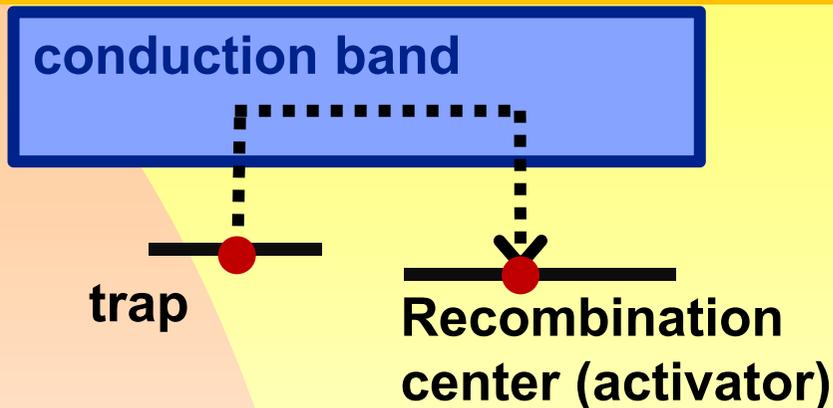
conduction band

trap

Recombination center (activator)

Escape from the trap

1. thermal escape via conduction band

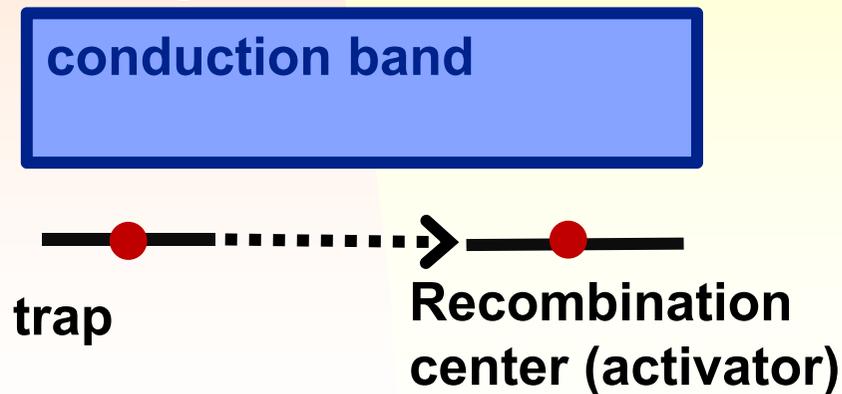


Shallow electron traps
Conduction band **is** involved

Similar scheme can apply for
holes and the valence band

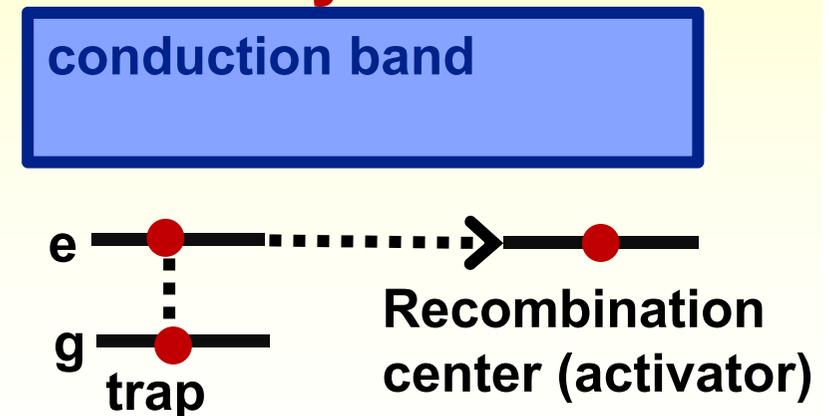
2. tunneling between traps and recombination centers

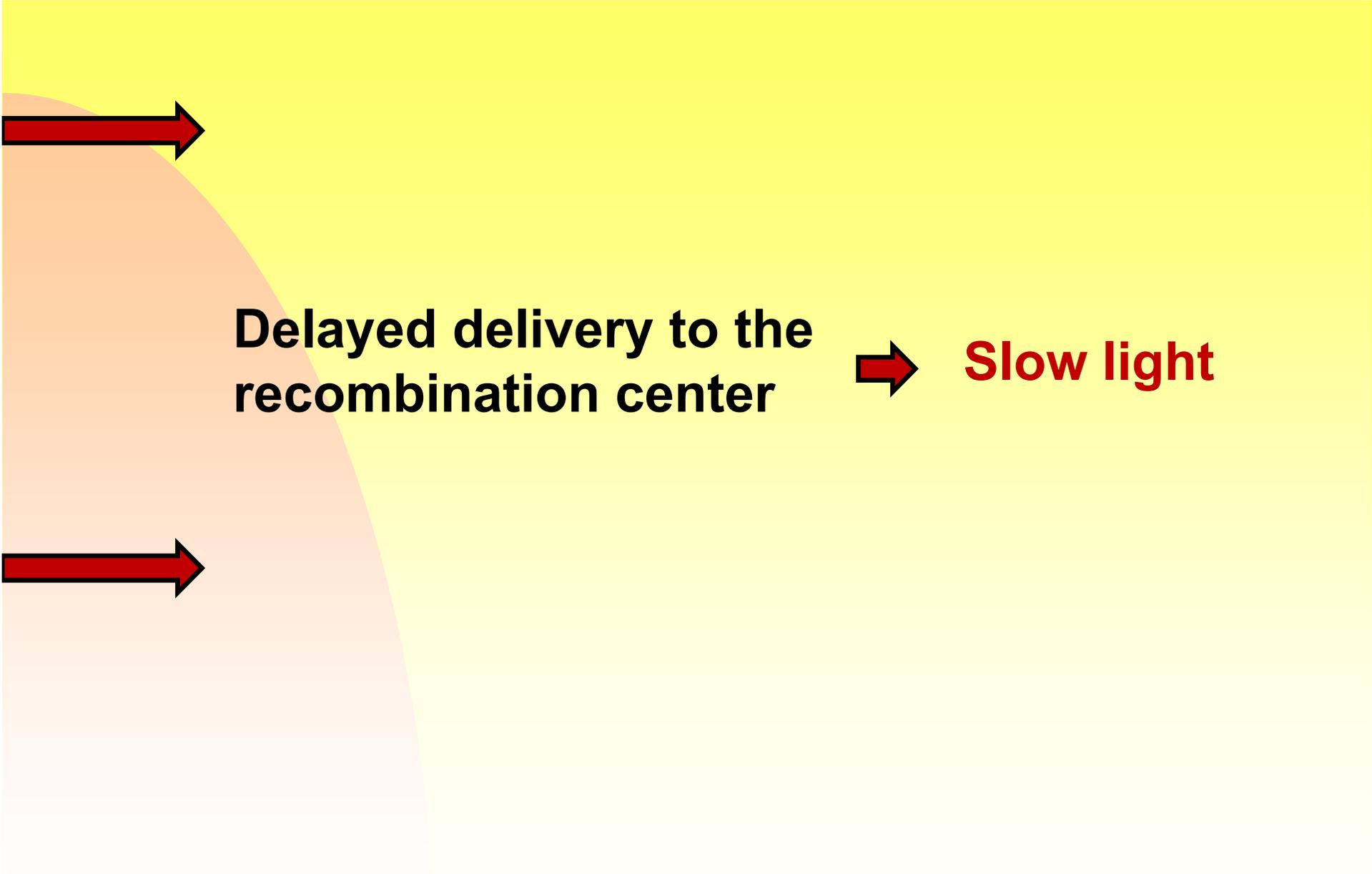
quantum



Shallow or deep electron traps
Conduction band **is not** involved

thermally assisted





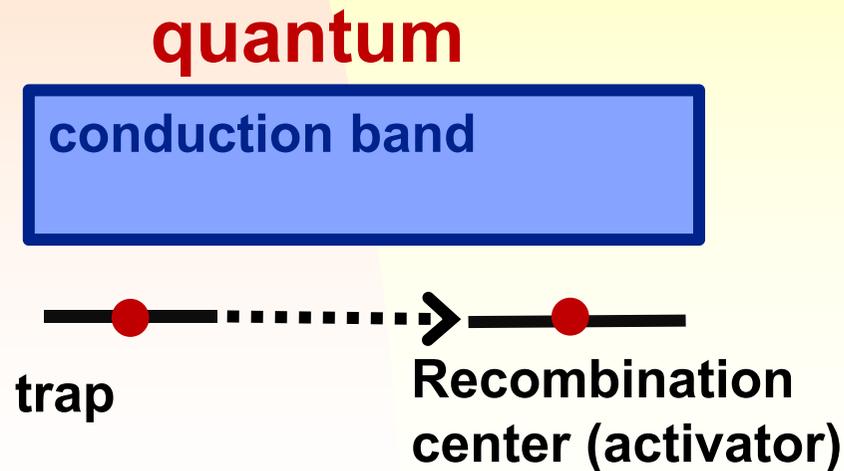
Delayed delivery to the recombination center



Slow light

Escape from the trap

2. tunneling between traps and recombination centers



Shallow or deep electron traps
Conduction band **is not** involved

How to detect the presence of quantum tunneling

Experimental tools

- monitoring the tail of the afterglow decay after X-ray irradiation
- scintillation decay measured over extended time scale
- thermally stimulated luminescence – glow peak analysis by the initial rise technique
- delayed recombination decay measurement

Tail of the afterglow decay

Quantum tunneling considered in the description of luminescence decay from early 1970s

C. J. Delbecq, Y. Toyozawa and P. H. Yuster, Phys. Rev. B **9**, 4497 (1974)

Anomalous fading of luminescence observed in materials used in TSL and OSL dating

From 1990s on



associated with quantum tunneling

Theoretical models of luminescence decay based on quantum tunneling between traps and **randomly distributed** recombination centers

$$I(t) \sim 1/t^p$$

$$p \sim 1$$

M. Tachyia and A. Mozumder, Chem. Phys. Lett. **28**, 87 (1974)

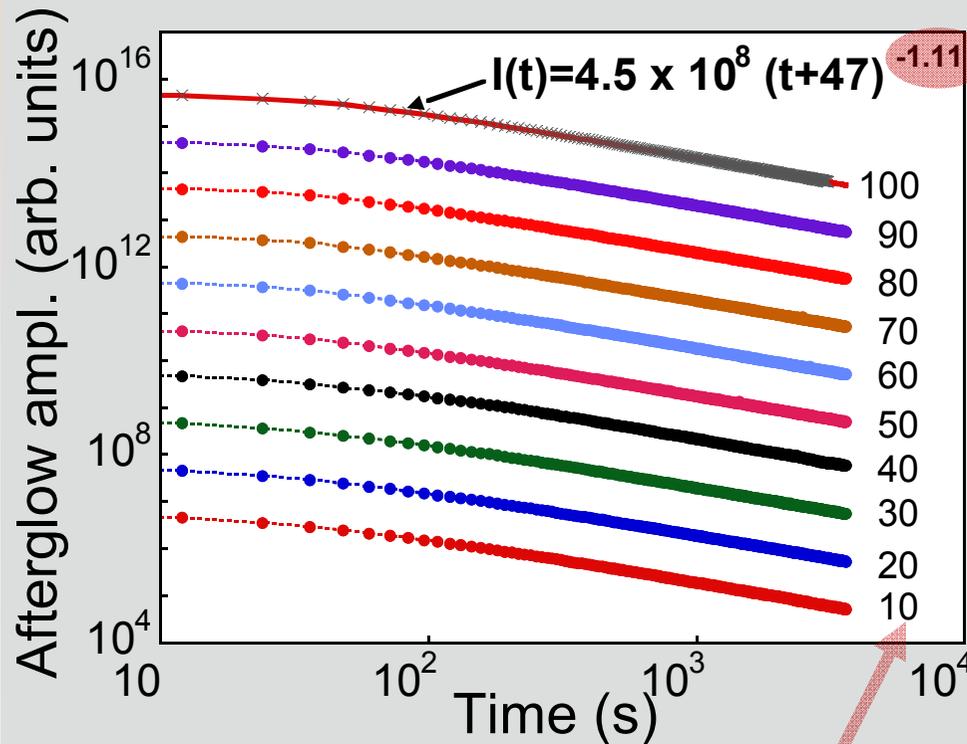
M. Tachyia and A. Mozumder, Chem. Phys. Lett. **34**, 77 (1975)

P. Avouris and T. N. Morgan, J. Chem. Phys. **74**, 4347 (1981)

D. J. Huntley, J. Phys.: Condens. Matter **18**, 1359 (2006)

Tail of the afterglow decay

LuAG:Ce³⁺



← Afterglow decay after X-ray irradiation at constant T

Fit formula:

$$I(t) = A \times (t + t_0)^{-p}$$

A, t_0 are constants

Exponent $p \sim 1$

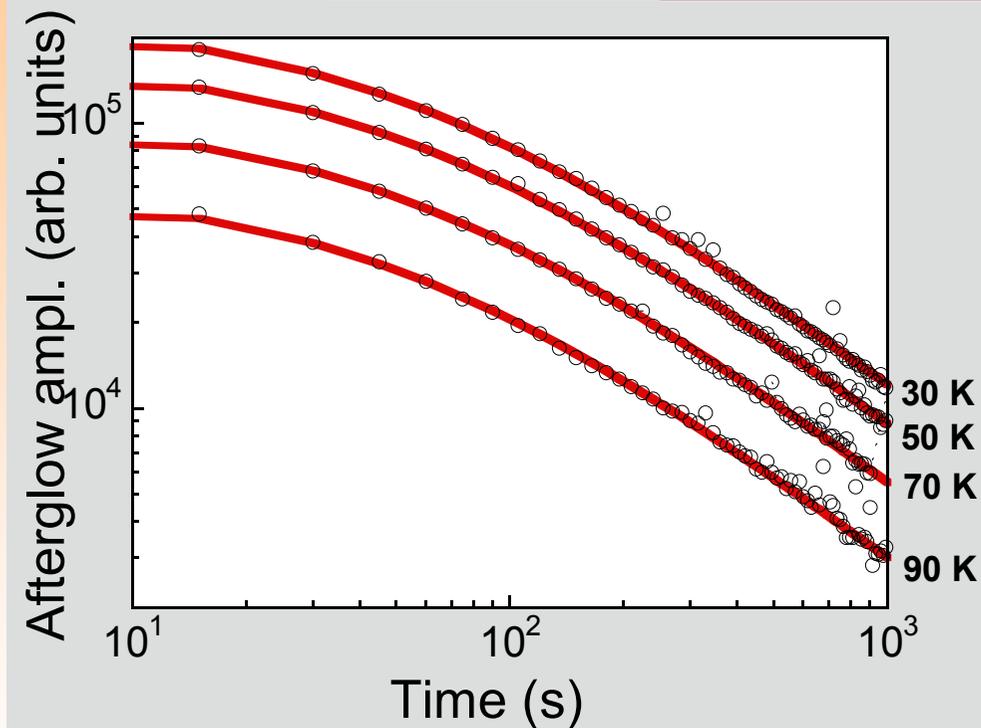
M. Nikl et al., Phys. Rev. B **76**, 195121 (2007)

Temperature in K

Decays are vertically shifted

↓
Afterglow tail is due to quantum tunneling of **electrons** between the traps and Ce³⁺ recombination centers – confirmed up to 100 K

Tail of the afterglow decay



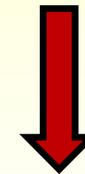
← Afterglow decay after X-ray irradiation at constant T

Decays for different T are vertically shifted

Fit formula:

$$I(t) = A/(t + t_0)$$

A, t_0 are constants

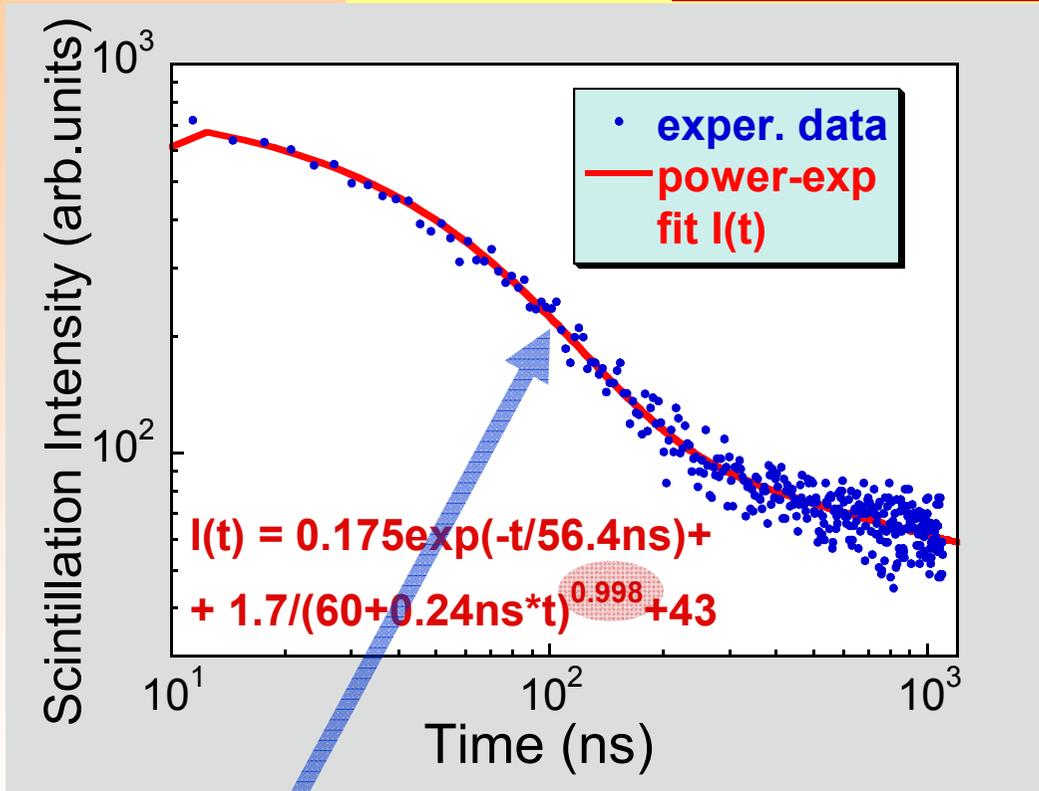


Afterglow tail is due to quantum tunneling of **holes** between the traps and recombination centers

E. Mihóková et al., Opt. Mater. **34**, 228 (2011)

Scintillation decay measured over extended time scale

LuAG:Ce³⁺



M. Nikl et al., Phys. Rev. B **76**, 195121 (2007)

← Scintillation decay at RT

Fit formula:

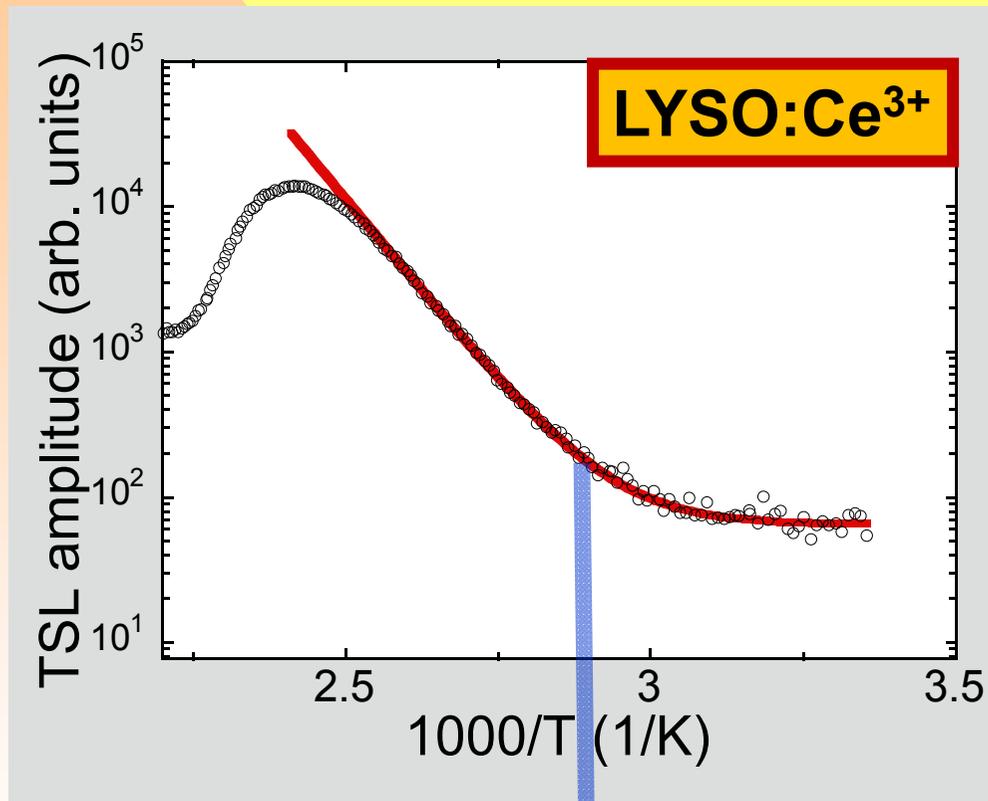
$$I(t) = B \times \exp(-t/\tau) + A/(t + t_0)^p + C$$

A, B, t_0 are constants,
 C is a background

Exponent $p \sim 1$

Linear part in log-log plot is due to quantum tunneling between the traps and recombination centers

TSL glow peak analysis by the initial rise technique



← Arrhenius plot after partial cleaning at 140 °C

frequency factor s

$$\beta E / k T_m^2 = s \times e^{-E/kT_m},$$

T_m - TSL peak maximum

β - heating rate

Detrapping time τ

$$\tau = s \times e^{E/kT_m},$$

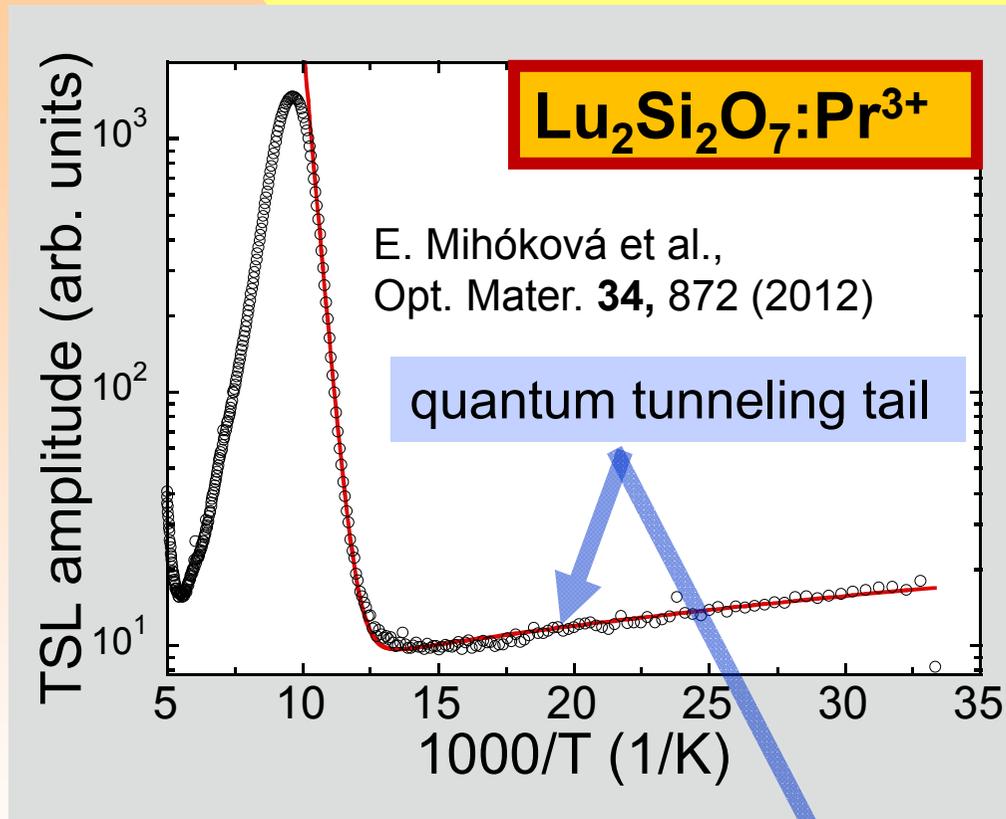
Fit function:

$$Amp(T) = b + w \times e^{-E/kT}$$

→ trap depth E

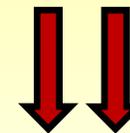
b - background, w - preexp. factor, k - Boltzmann constant, A - constant

TSL glow peak analysis by the initial rise technique



← Arrhenius plot after partial cleaning at 90 K

quantum tunneling tail $\sim 1/t$
linear heating $T \sim t$



quantum tunneling tail $\sim 1/T$

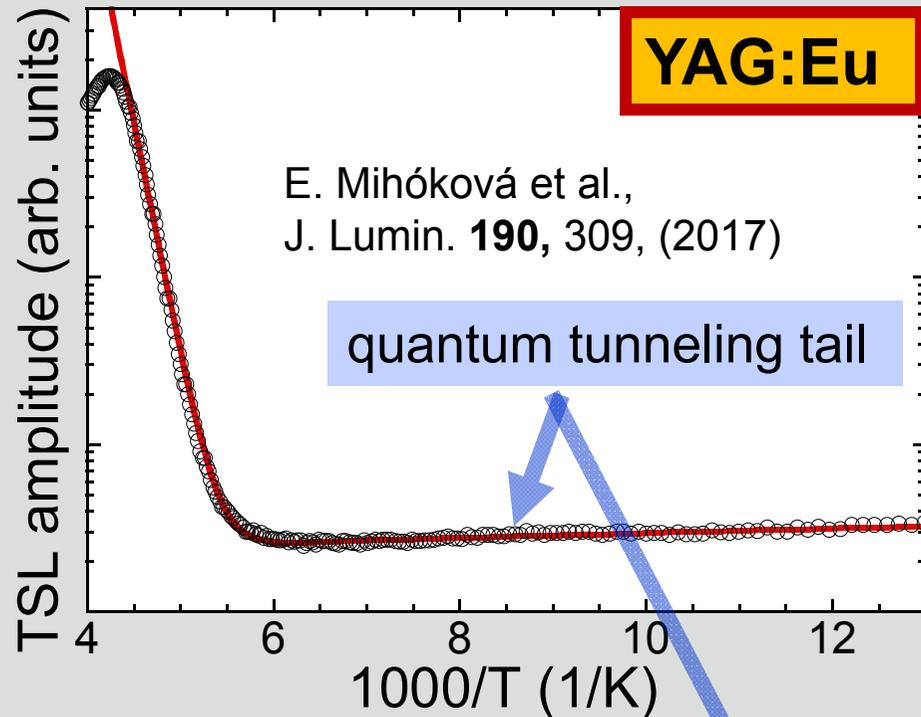
Fit function:

$$Amp(T) = b + w \times e^{-E/kT} + \frac{A}{T}$$

→ trap depth E

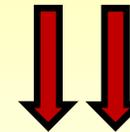
b - background, w - preexp. factor, k - Boltzmann constant, A - constant

TSL glow peak analysis by the initial rise technique



← Arrhenius plot after partial cleaning at 219 K

quantum tunneling tail $\sim 1/t$
linear heating $T \sim t$



quantum tunneling tail $\sim 1/T$

Fit function:

$$Amp(T) = b + w \times e^{-E/kT} + \frac{A}{T}$$

→ trap depth E

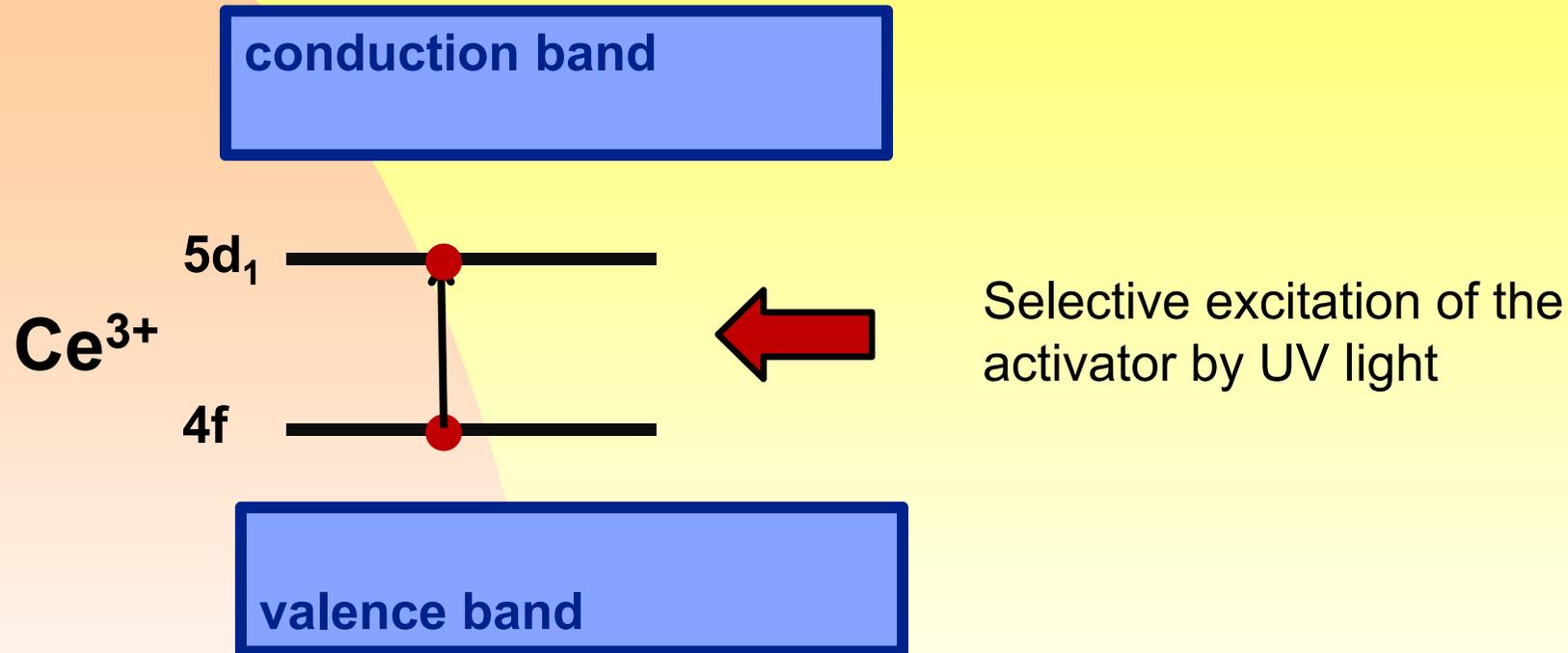
b - background, w - preexp. factor, k - Boltzmann constant, A - constant

Delayed recombination decay

Originally developed as a contactless purely optical technique to study the process of thermal ionization of the activator's excited state

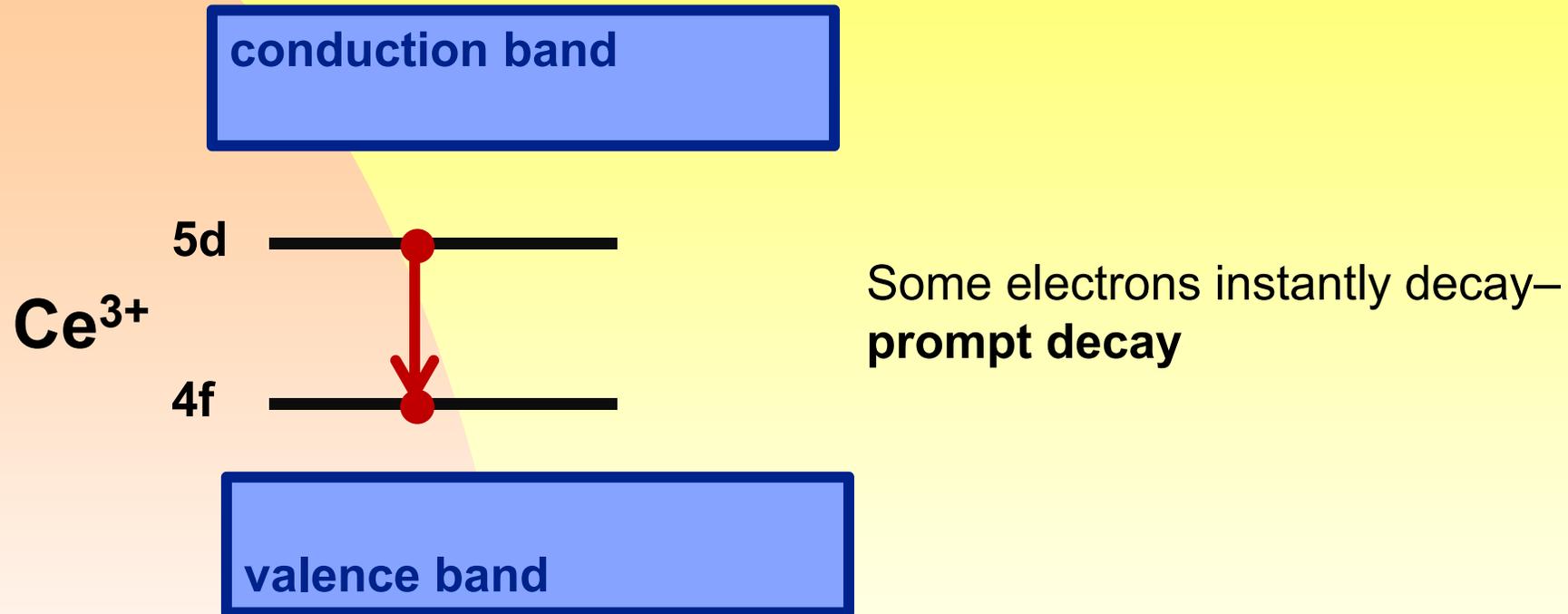
Delayed recombination decay

Principle



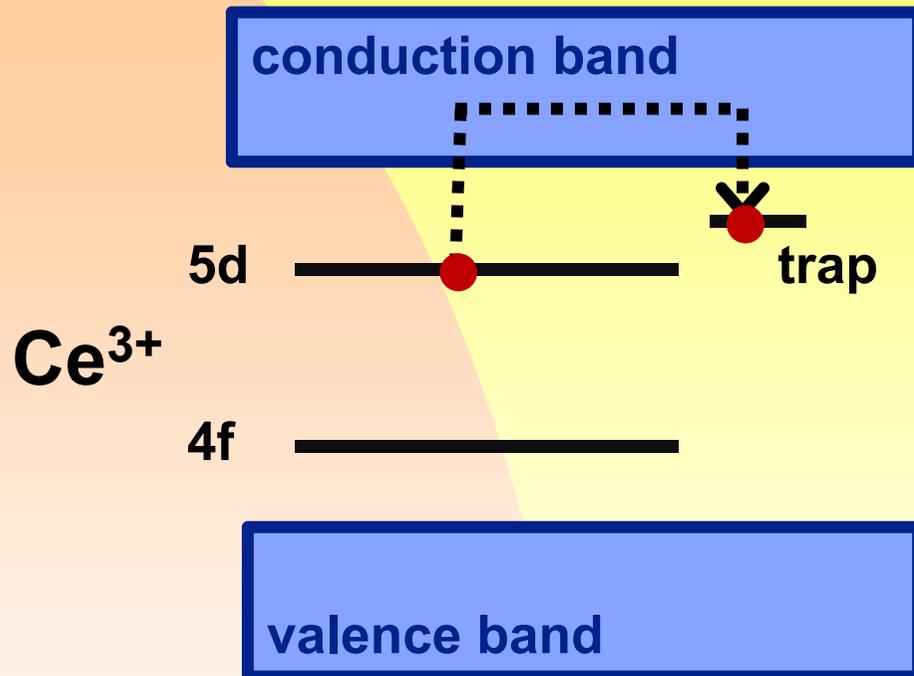
Delayed recombination decay

Principle



Delayed recombination decay

Principle

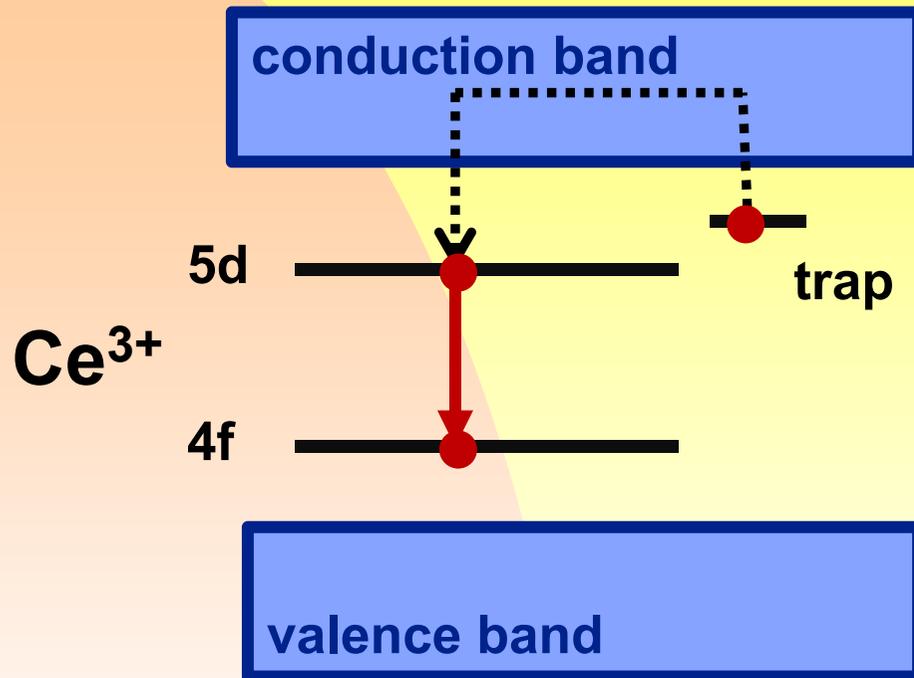


Some electrons **do not** decay promptly

If there is sufficient thermal energy around, electron can escape from the excited state into the conduction band and be trapped

Delayed recombination decay

Principle



Electron **after some time** spent in the trap escapes and can be recaptured at the luminescence center

Electron recombines with the hole in the ground state and produces the **delayed recombination light**



slow tail in the luminescence decay

Delayed recombination decay vs **afterglow**

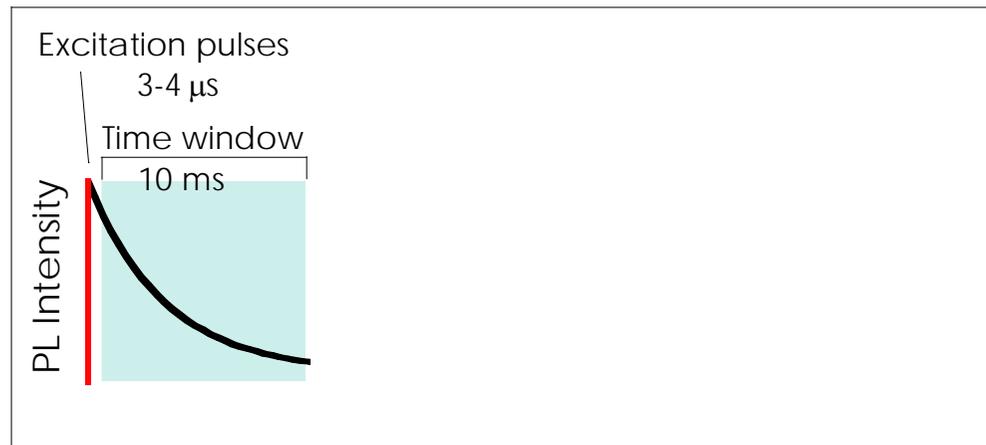
Selective UV excitation

excitation by ionizing radiation

Monitoring the T dependence of the decay supplies the information about the excited state thermal ionization

Delayed recombination decay

Measurement



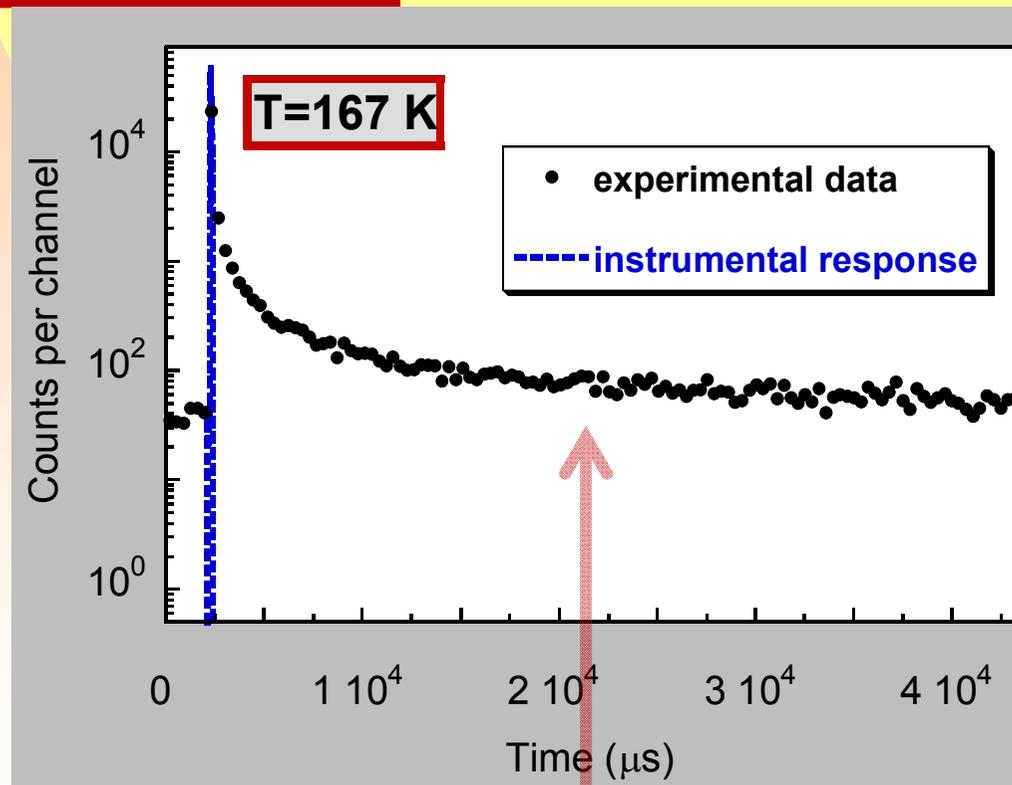
PL decay measurement

the signal is accumulated in the **extended** time window (up to 10 minutes) for a set of different temperatures

Delayed recombination decay

Example

SrHfO₃:Ce³⁺



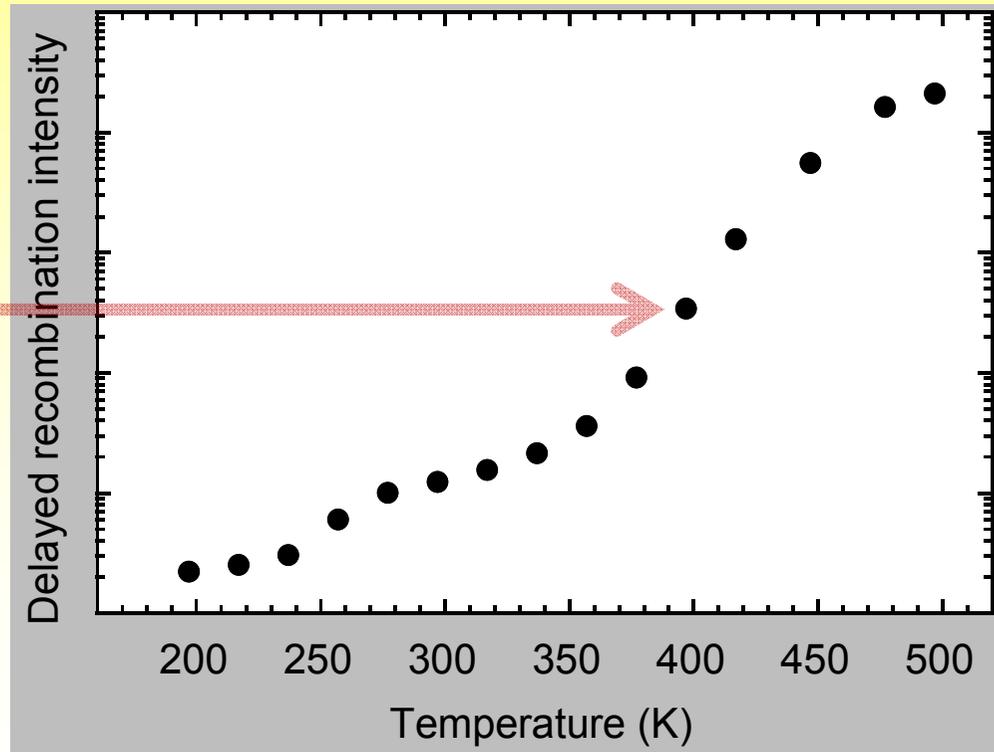
For each temperature we integrate the **decay** in time within the time (interval) of the measurement

Temperature dependence of delayed recombination decay

Experiment



Each data point is time integrated decay at the temperature T



Temperature dependence of delayed recombination decay

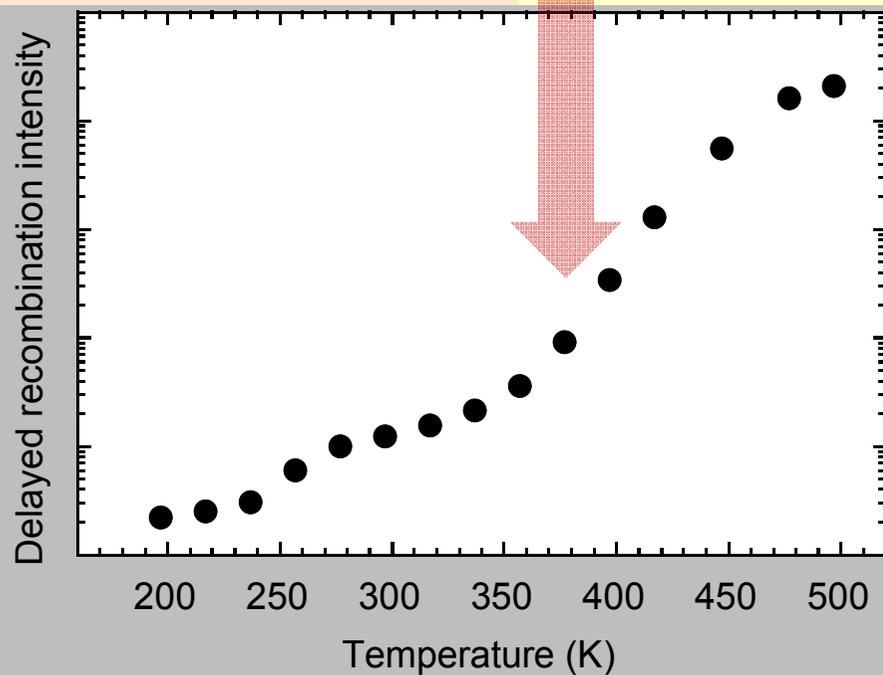
Model

$$I_{DR}(T) = w e^{-\frac{E_{ion}}{kT}} \sum_{k=1}^n A_k \left[e^{-\frac{t_b}{\tau(E_k)}} - e^{-\frac{t_e}{\tau(E_k)}} \right]$$

M. Fasoli, A. Vedda, E. Mihóková, M. Nikl, PRB 85, 085127 (2012)

Lu₂Si₂O₇:Pr³⁺

crystal



Temperature dependence of delayed recombination decay

Model

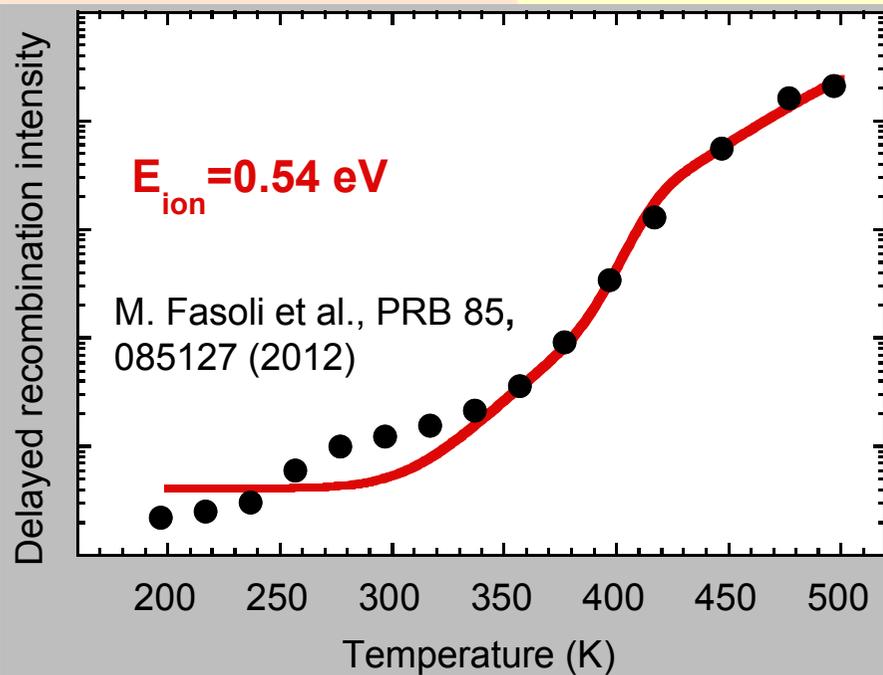
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M. Fasoli, A. Vedda, E. Mihóková, M. Nikl, PRB 85, 085127 (2012)

Ionization energy is determined from the fit

Lu₂Si₂O₇:Pr³⁺

crystal



Temperature dependence of delayed recombination decay

Model

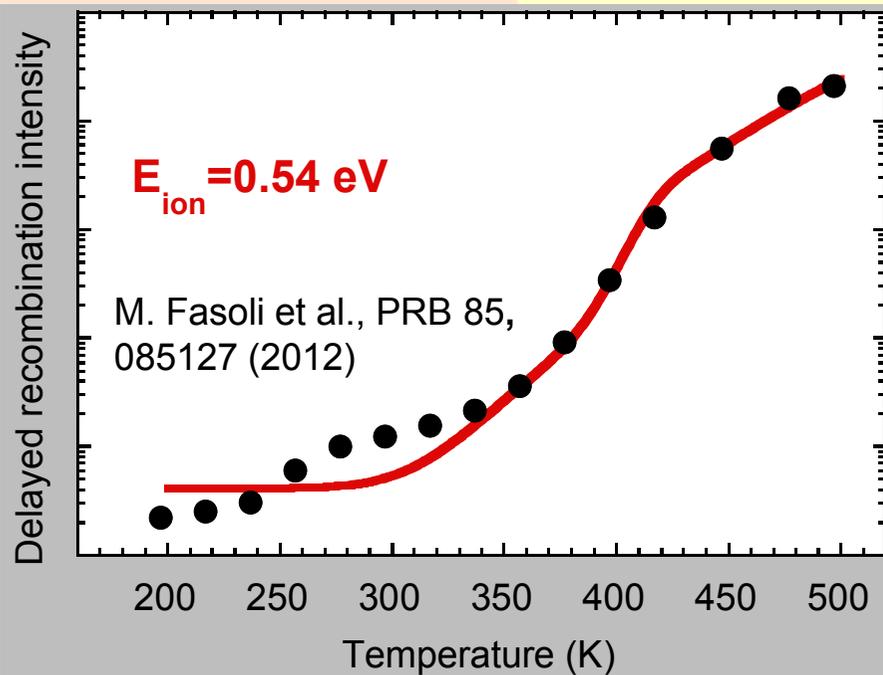
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M. Fasoli, A. Vedda, E. Mihóková, M. Nikl, PRB 85, 085127 (2012)

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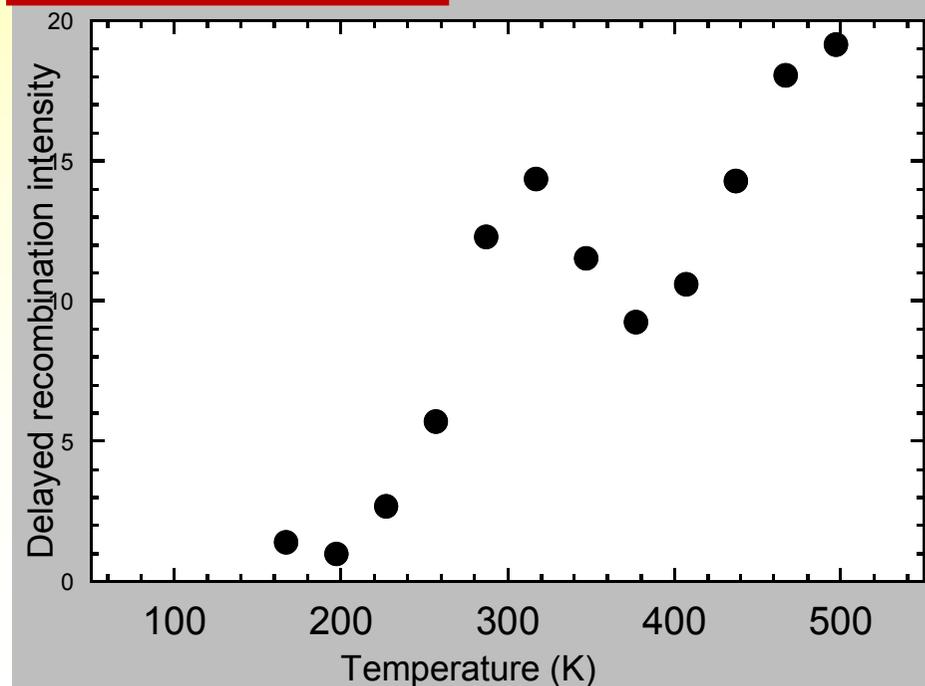
Lu₂Si₂O₇:Pr³⁺

crystal



SrHfO₃:Ce³⁺

powder



Temperature dependence of delayed recombination decay

Model

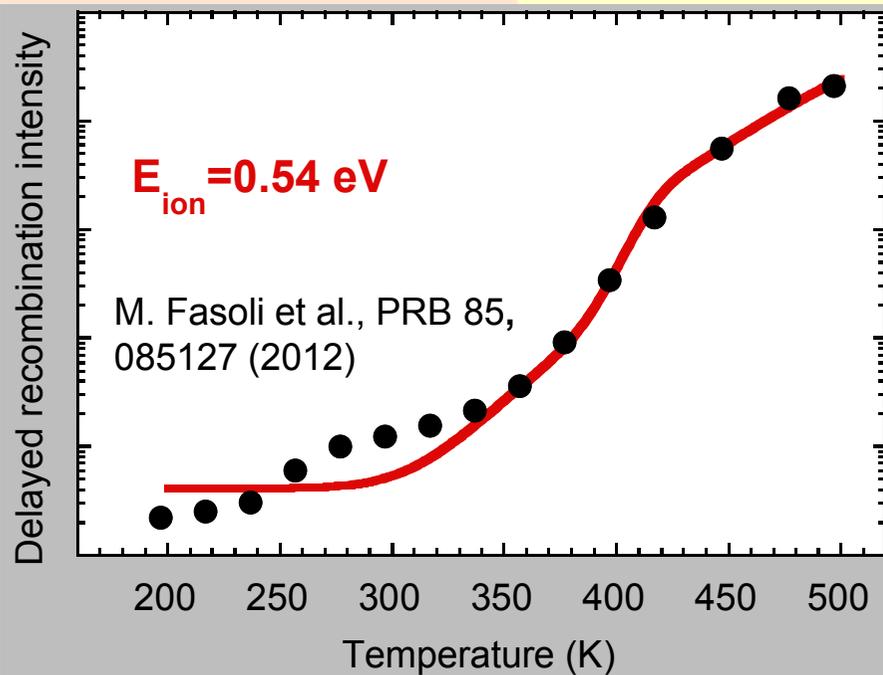
$$I_{DR}(T) = w e^{-\frac{E_{ion}}{kT}} \sum_{k=1}^n A_k \left[e^{-\frac{t_b}{\tau(E_k)}} - e^{-\frac{t_e}{\tau(E_k)}} \right]$$

M. Fasoli, A. Vedda, E. Mihóková, M. Nikl, PRB 85, 085127 (2012)

Ionization energy is determined from the fit

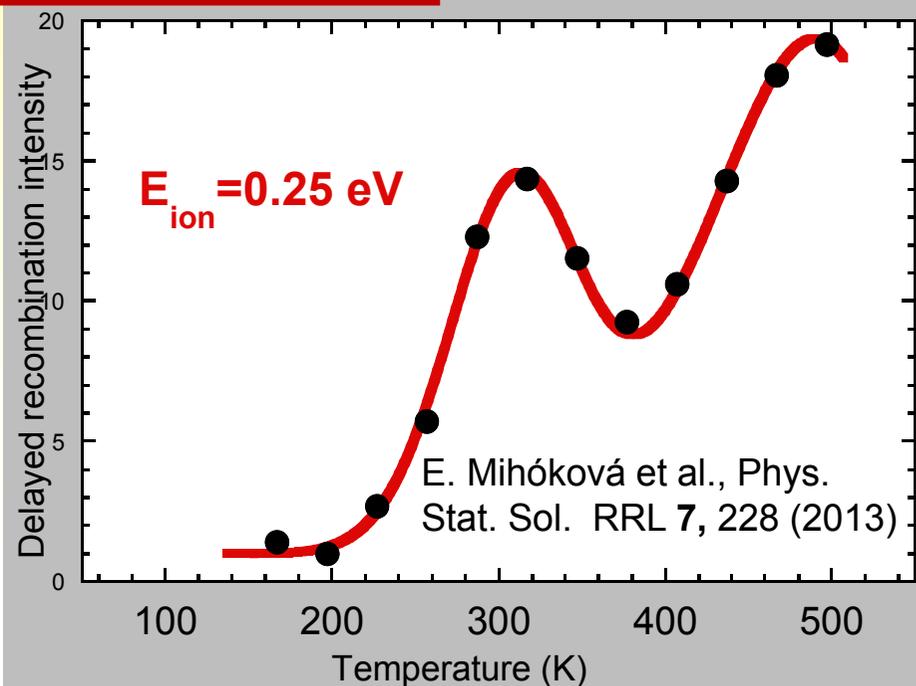
Lu₂Si₂O₇:Pr³⁺

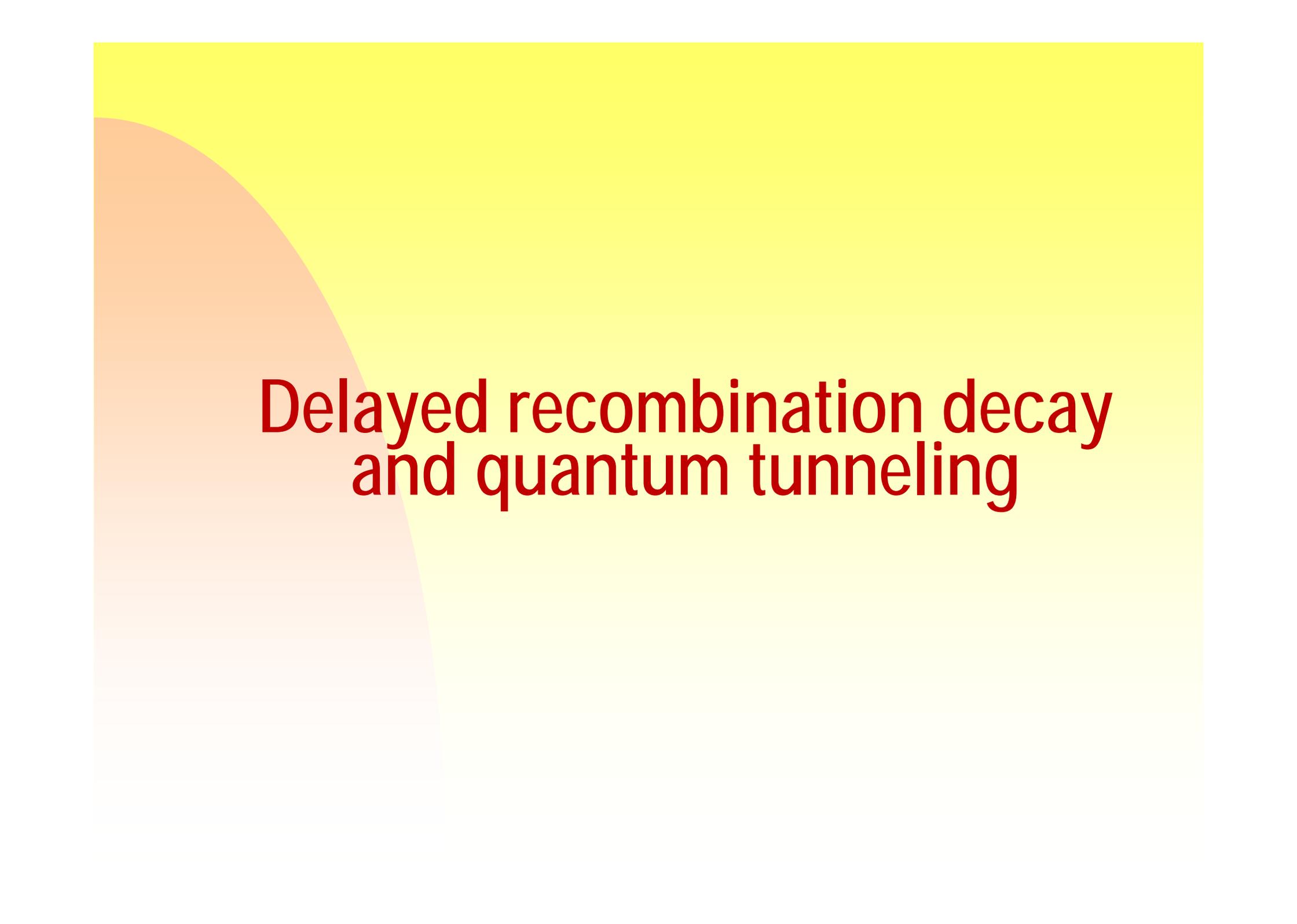
crystal



SrHfO₃:Ce³⁺

powder

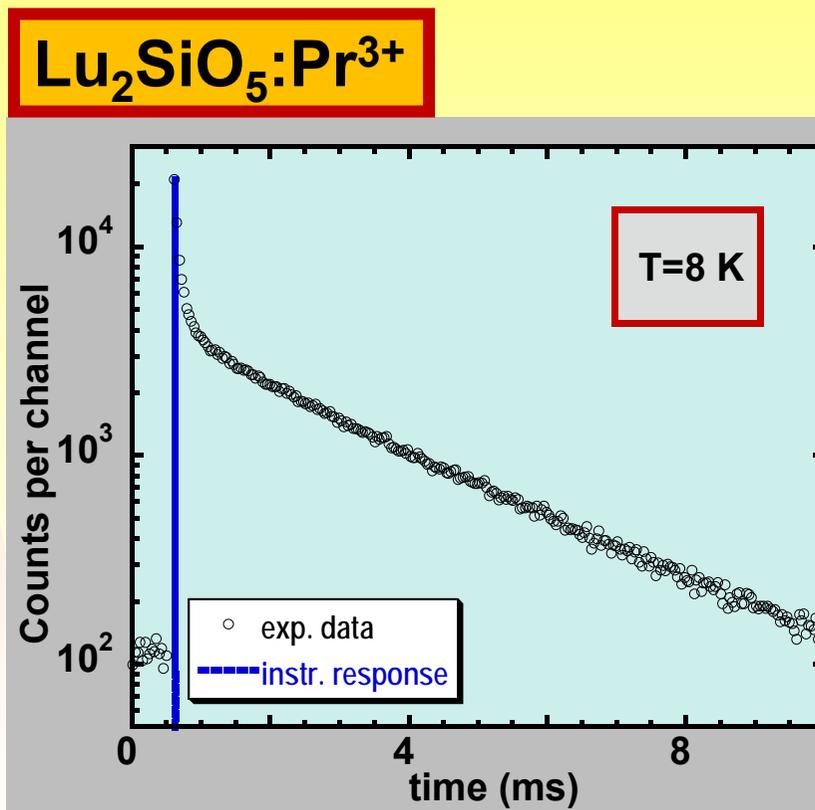




Delayed recombination decay and quantum tunneling

Low temperature contribution to delayed recombination decay

Experiment

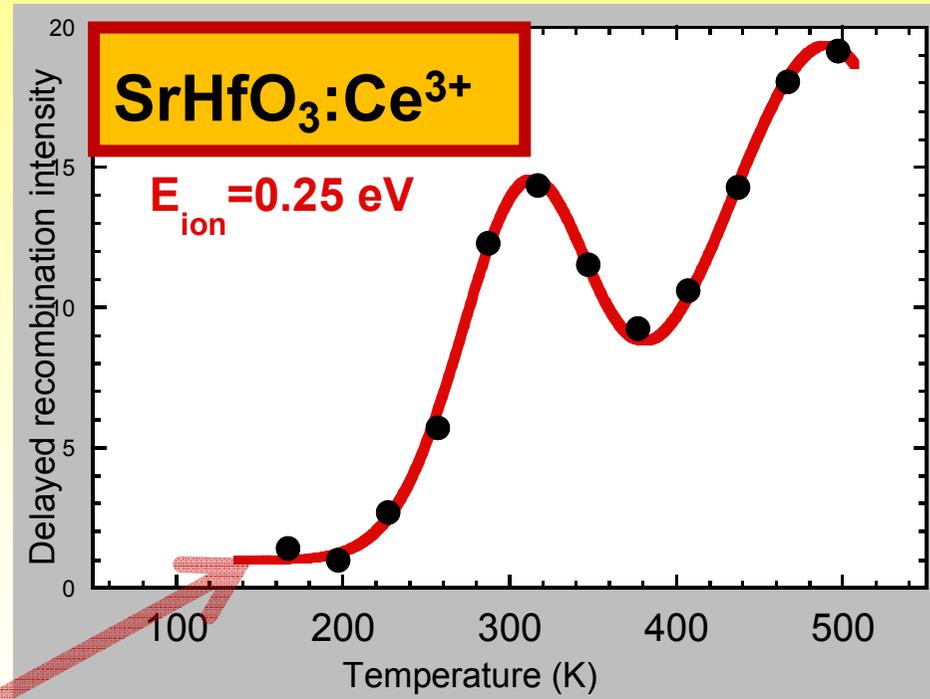
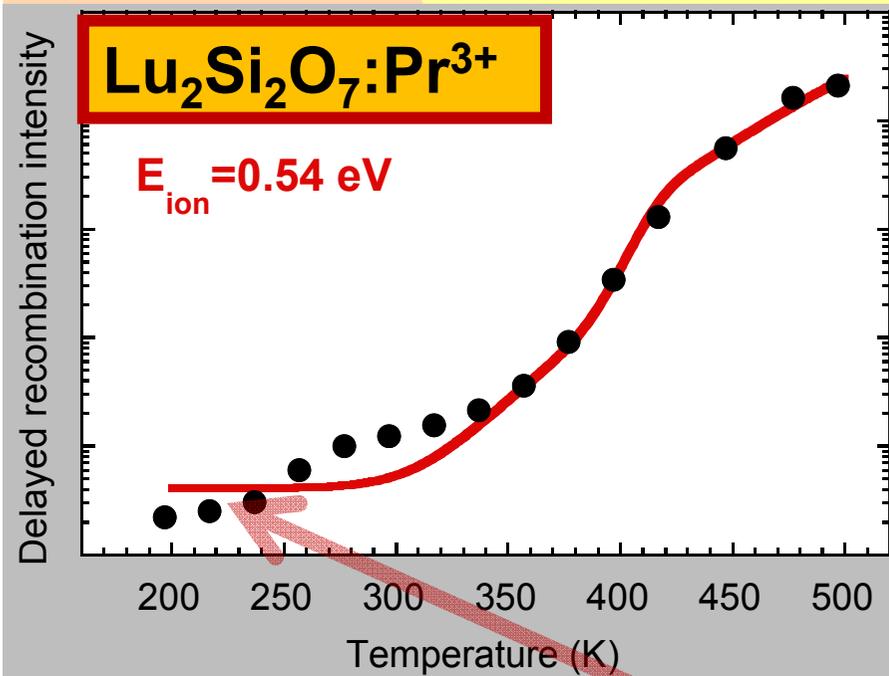


There is **low T contribution** to the delayed recombination that is **NOT** due to thermal ionization

Fit of the model to the data

$$I_{DR}(T) = w e^{-\frac{E_{ion}}{kT}} \sum_{k=1}^n A_k \left[e^{-\frac{t_b}{\tau(E_k)}} - e^{-\frac{t_e}{\tau(E_k)}} \right]$$

Low T limit of the modelled $I_{DR}(T)$ is zero



Actual fit requires an additive constant to reproduce the low T data



Origin of the low temperature contribution



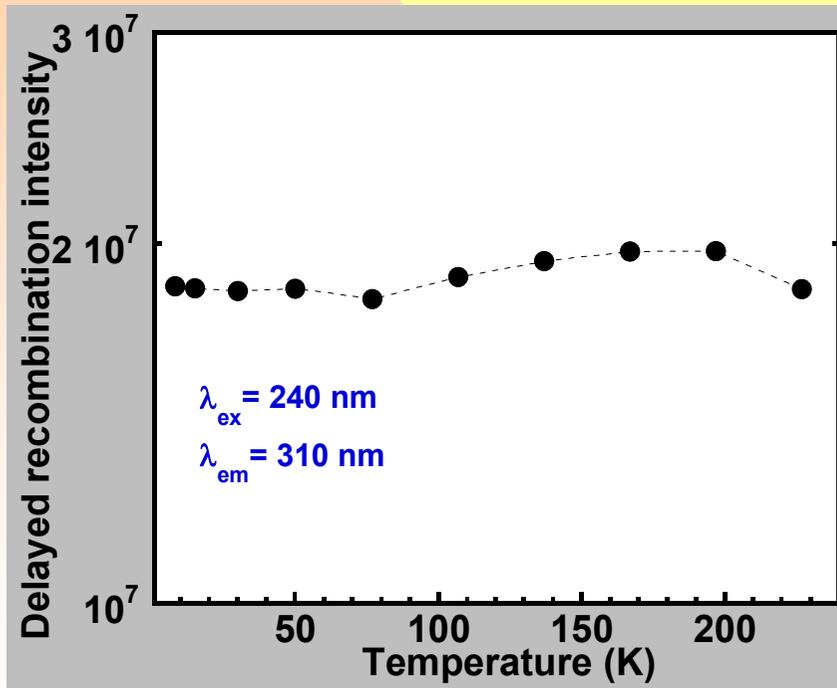
quantum tunneling
between the activator and trap(s)

Experimental verification: observation of temperature independent delayed recombination intensity within certain low T interval

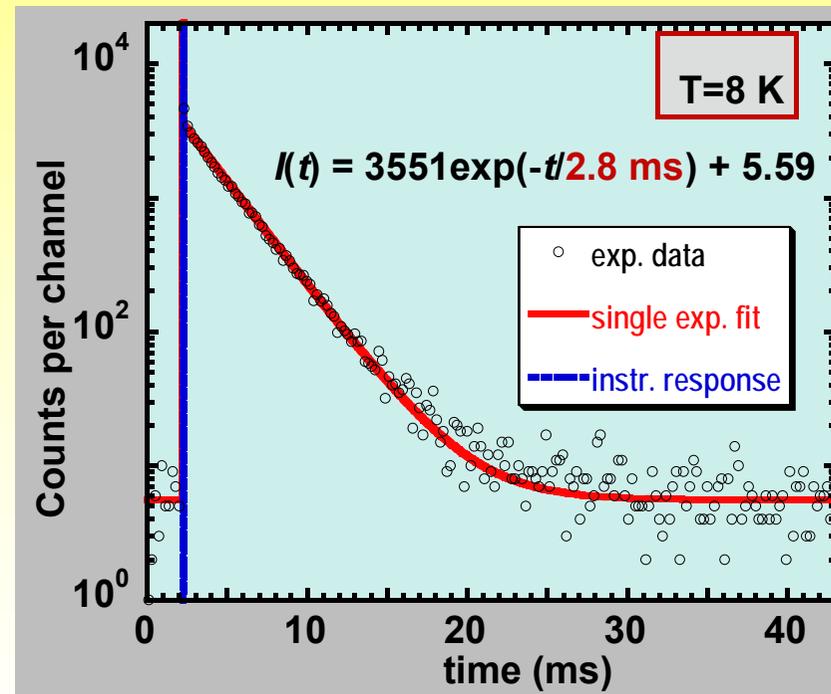
Theoretical verification: modelling of quantum tunneling shows it can cause several orders of magnitude delay of the decay

Experimental support

Delayed recombination decay at the low temperature limit



constant signal up to about 200 K

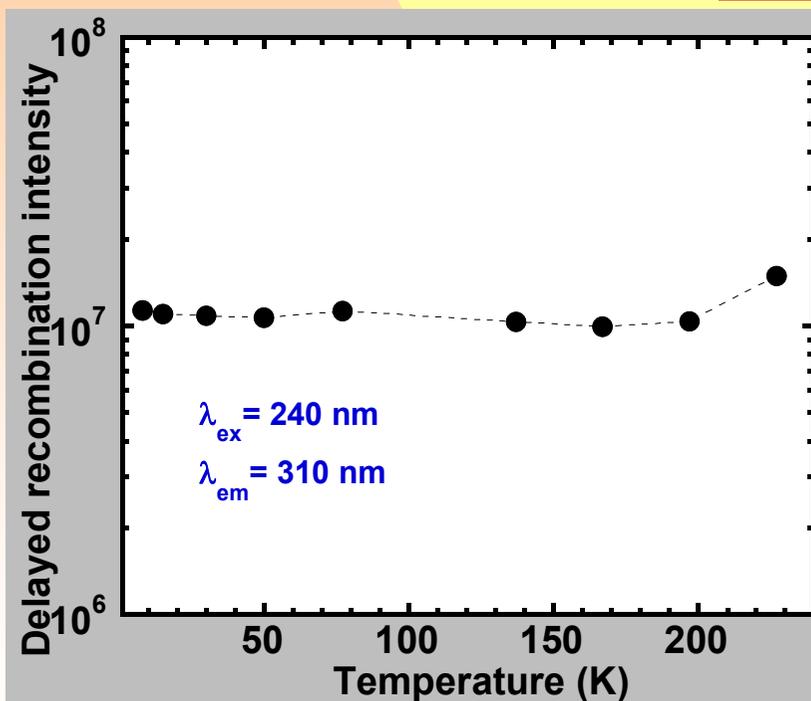


single exponential decay up to 200 K

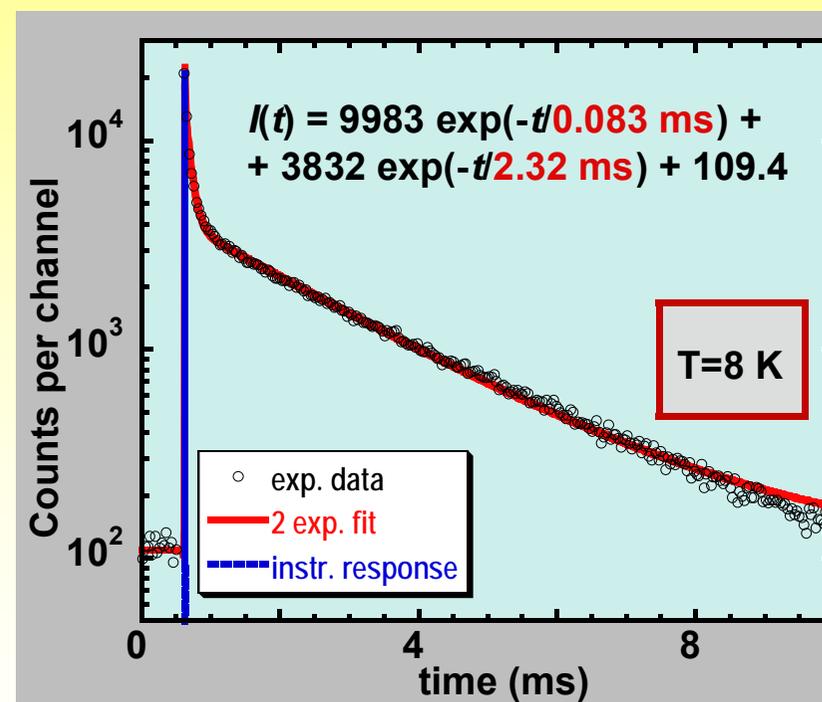
Quantum tunneling between the activator and a single trap

Experimental support

Delayed recombination decay at the low temperature limit



constant signal up to about 200 K



complex decay

Quantum tunneling between the activator and several traps

Experimental support

Delayed recombination decay at the low temperature limit

Quantum tunneling experimentally verified in a variety of Ce^{3+} and Pr^{3+} -doped scintillating complex oxides in various forms (single crystals, thin films, nanopowders)

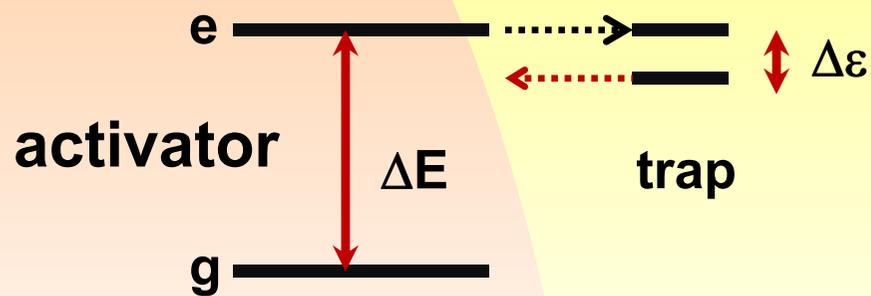
E. Mihóková et al., IEEE Trans. Nucl. Sci. **61**, 257 (2014)

E. Mihóková et al., Opt. Mater. **40**, 127 (2015)

Theoretical support

One dimensional model of quantum tunneling between the trap and activator

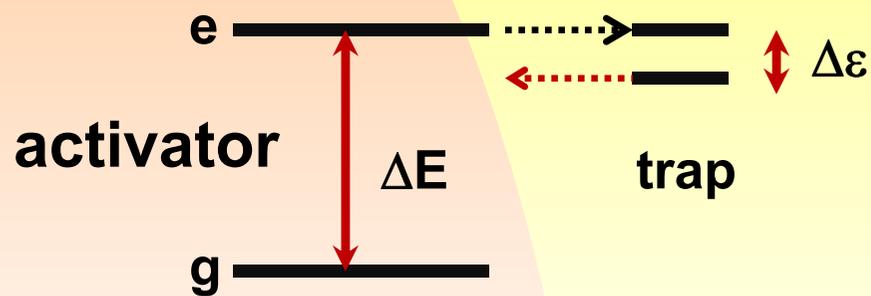
Scenario



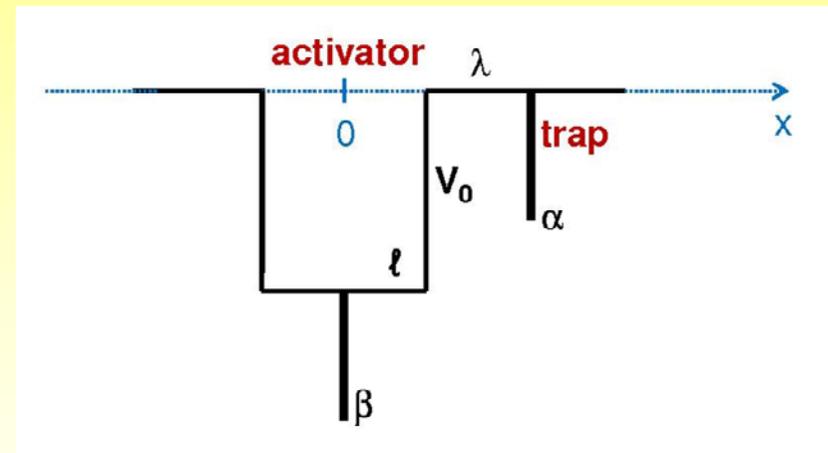
Theoretical support

One dimensional model of quantum tunneling between the trap and activator

Scenario



Model potential



$$V = -\beta\delta(x) - V_0 \theta(\ell - |x|) - \alpha\delta(x - \ell - \lambda)$$

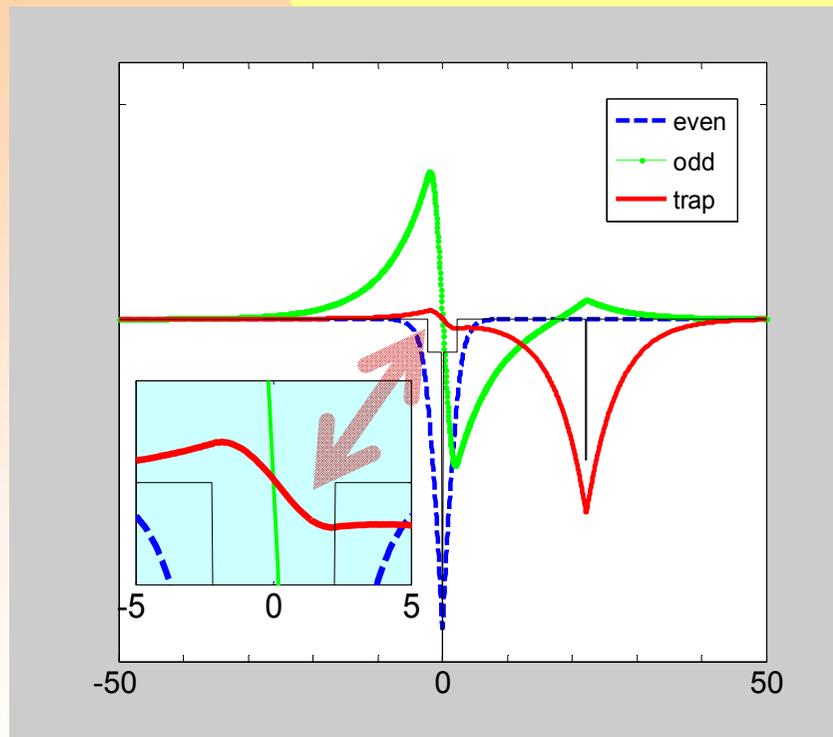
We solved the Schrödinger equation and calculated the **decay rate slowed by quantum tunneling**

$$\Gamma = R \times \Gamma_0$$

slow rate reduction factor prompt rate

Theoretical support

One dimensional model of quantum tunneling between the trap and activator



Wave functions

even= ground state of the activator

odd= excited state of the activator

trap= trap state

trap wave function in the activator region

→ its norm gives the reduction factor **R**

With reasonable parameter values the decay is delayed by **several orders of magnitude** which is what we see in the experiment

Theoretical support

Two dimensional model of quantum tunneling between the trap and activator

Trap and activator – isotropic Gaussian wells

Model potential

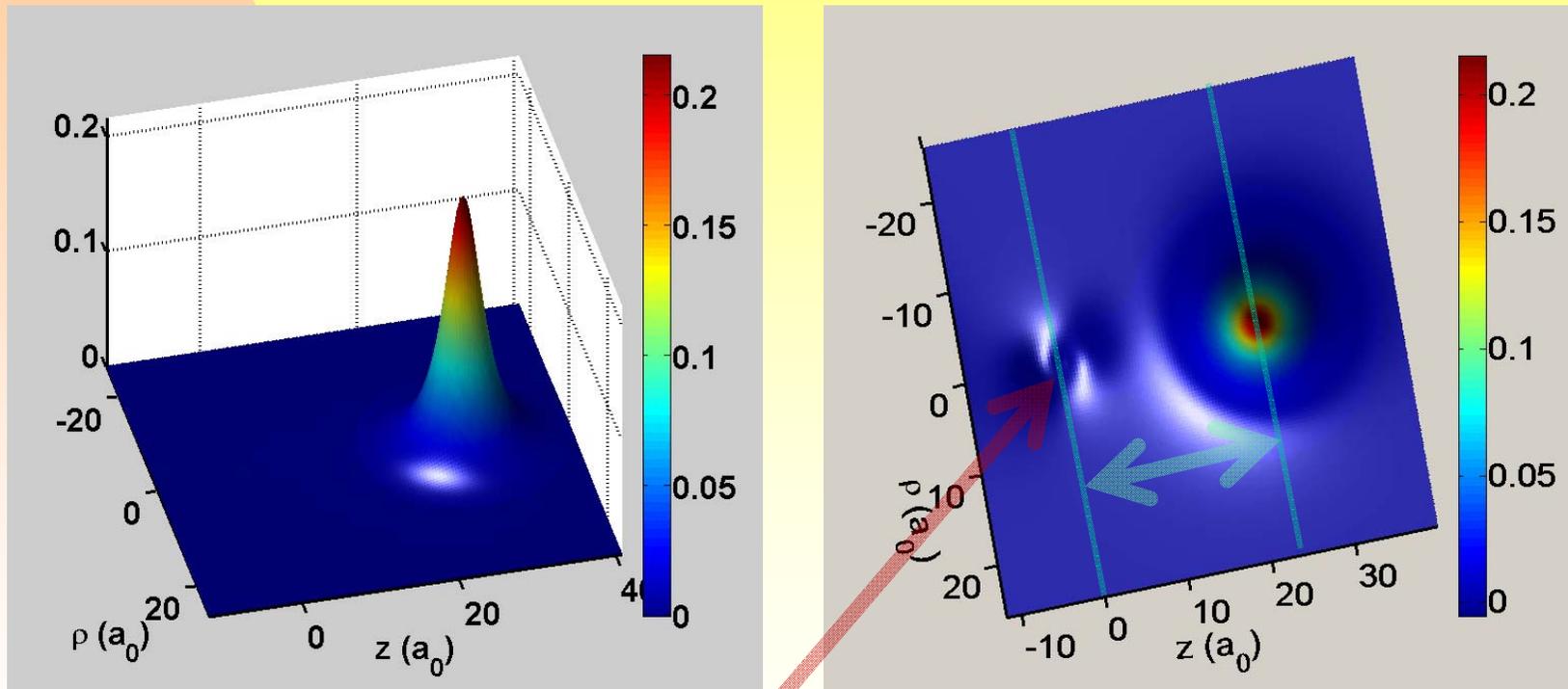
$$V(\rho, z) = -\alpha e^{-(\rho^2 + z^2)/r_A^2} + \frac{\beta}{2(\rho^2 + z^2)} - \gamma e^{-(\rho^2 + (z-\lambda)^2)/r_T^2}$$

Solution of the Schrödinger equation by the “Wave packet” program

E. Mihóková and L. S. Schulman, J. Phys. Cond. Matter **27**, 075501 (2015)

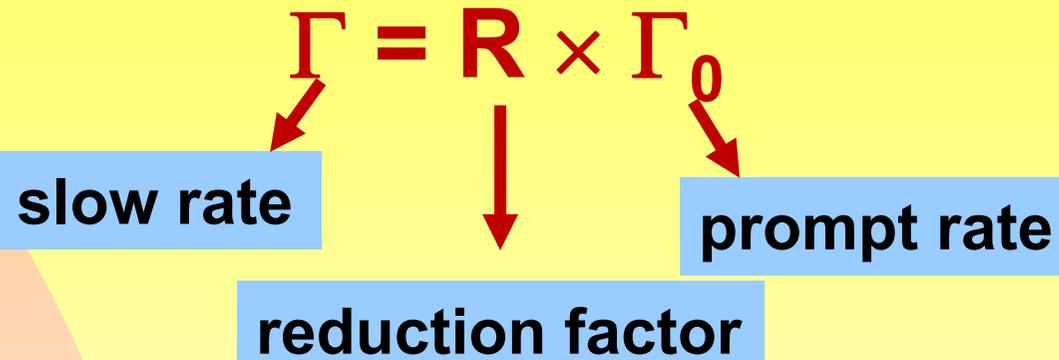
Theoretical support

Two dimensional model of quantum tunneling between the trap and activator



trap wave function shown from two different perspectives
the portion of the wave function in the activator region
trap – activator distance is $23 a_0 \sim 12 \text{ \AA}$
the delay of the decay ~ 5 orders of magnitude

Analytical estimates of R



ground state: E_G, ψ_G
 excited state: E_E, ψ_E
 trap state: E_T, ψ_T

$$\Gamma_0 \sim |\langle \psi_E | r | \psi_G \rangle|^2$$

trap wave function dropoff in 3D:

$$\sqrt{\kappa/2\pi} \exp(-\kappa r)/r$$

r - distance from the trap center

$$\kappa = \sqrt{-2mE_T/\hbar^2}$$

Because in the region of the activator ψ_T resembles ψ_E , the reduction, R , arises simply from the smaller amplitude of ψ_T

$$R = \frac{\eta \int_{\text{activator}} \left| \sqrt{\frac{\kappa}{2\pi}} \frac{\exp(-\kappa r)}{r} \right|^2 d^3 r_{\text{activator}}}{\langle \psi_T | \psi_T \rangle}$$

$\eta = 2-3$, correction factor

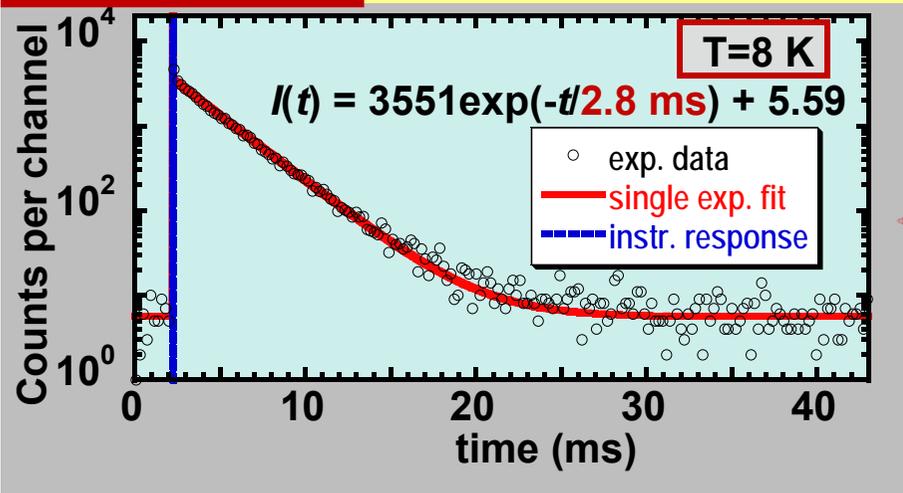
Application

$\text{Y}_2\text{SiO}_5:\text{Pr}^{3+}$

single exp. decay



Quantum tunneling between the activator and a single (kind of) trap



$$1/\Gamma_0 = 18 \text{ ns}$$

$$1/\Gamma = 2.8 \text{ ms}$$

$$R = 8.2 \times 10^{-6}$$

$$\text{trap depth } E_T \approx E_E \sim 0.4 \text{ eV}$$

J. Pejchal et al., J. Phys. D:
Appl. Phys. **42** 055117 (2009)

J. Pejchal et al., J. Phys. D:
Appl. Phys. **42** 055117 (2009)

delayed recombination decay

Analytical estimates of R + experimental data



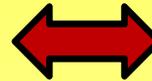
Trap \leftrightarrow Pr^{3+} distance

$$\lambda = 12\text{-}14 \text{ \AA}$$

Application

Lu₂Si₂O₇:Pr³⁺

complex decay



Quantum tunneling between the activator and traps at various distances from the activator

M. Nikl et al., Chem. Phys. Lett. **493** 72 (2010)

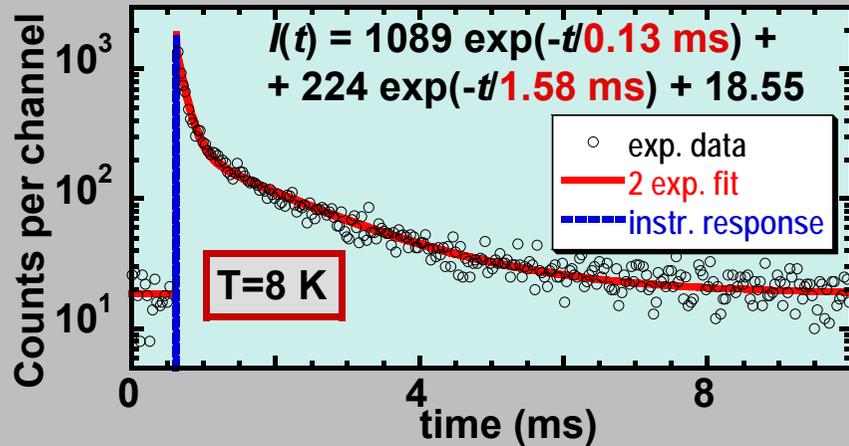
$$1/\Gamma_0 = 19 \text{ ns}$$

$$1/\Gamma_1 = 0.13 \text{ ms} \quad 1/\Gamma_2 = 1.58 \text{ ms}$$

$$R_1 = 1.5 \times 10^{-4} \quad R_2 = 1.3 \times 10^{-5}$$

$$\text{trap depth } E_T \approx E_E = 0.54 \pm 0.05 \text{ eV}$$

M. Fasoli et al., Phys. Rev. B **85**, 085127 (2012)



delayed recombination decay

Analytical estimates of R + experimental data



Trap ↔ Pr³⁺ distance

$$\lambda_1 = 6-8 \text{ \AA}$$

$$\lambda_2 = 10-11 \text{ \AA}$$

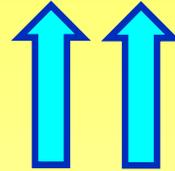
Independent TSL experiment confirms the presence of the trap with $E_T = 0.51 \pm 0.01 \text{ eV}$

E. Mihóková et al., Opt. Mater. **34**, 872 (2012)

Trap is likely to be an oxygen vacancy

Mean distance Lu-O in Lu₂Si₂O₇:Pr³⁺ is 2.23 Å → 3rd or 4th nearest O site

Due to quantum tunneling between the activator and traps



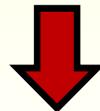
accumulated experimental and theoretical support



Low T delayed recombination decay



An experimental tool to detect the presence of quantum tunneling in studied scintillating materials



With analytical estimates and additional information from independent experiments provides a handle on trap-activator distances

Summary

- **Quantum tunnelling between traps and recombination centers belongs to sources of losses of fast scintillation light**
- **Quantum tunneling is quite frequently encountered in scintillating materials**
- **Detection of the presence of quantum tunnelling:**
 - monitoring the afterglow tail**
 - scintillation decay**
 - initial rise in TSL**
 - delayed recombination decay**



Thanks are expressed to all collaborators

and to

the audience for kind attention